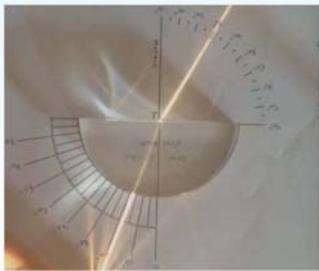
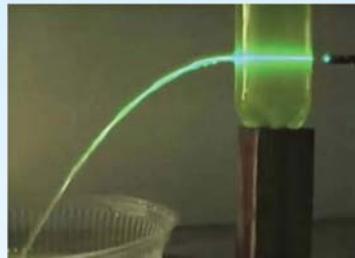
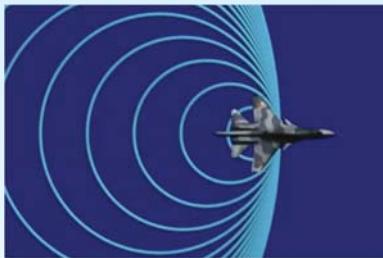




Vågrörelselära och optik



Kapitel 16 - Ljud

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1



Vågrörelselära och optik



Kurslitteratur: University Physics by Young & Friedman

Harmonisk oscillator:

Kapitel 14.1 - 14.4

Mekaniska vågor:

Kapitel 15.1 - 15.8

Ljud och hörande:

Kapitel 16.1 - 16.9

Elektromagnetiska vågor:

Kapitel 32.1 & 32.3 & 32.4

Ljusets natur:

Kapitel 33.1 - 33.4 & 33.7

Stråloptik:

Kapitel 34.1 - 34.8

Interferens:

Kapitel 35.1 - 35.5

Diffraktion:

Kapitel 36.1 - 36.5 & 36.7



Vågrörelselära och optik



Tid	Må	02-nov	Ti	03-nov	On	04-nov	To	05-nov	Fr	06-nov
08-10	Kvantfysik (A)	Våglära/optik (A)	kap 14	Kvantfysik (A)	Våglära/optik (A)	Våglära/optik (A)	Våglära/optik (A)	Kvantfysik (A)	Kvantfysik (A)	
10-12	Intro period 2 (A) Informationssökning (A)	Kvantfysik (A)	Kvantfysik kap 14+15	Våglära/optik (A)	ÄFYA11 (L218)	Kvantfysik (A)	Kvantfysik (A)	Kvantfysik (A)	Kvantfysik (A)	kap 15
13-15	Utvärdering (A) 12-13:	Övningar Optik&Våg (L218-19)	Sl gp6-10 (L219)	ÄFYA11 (L218)	Sl gp11-15 (L219)				Övningar Optik&Våg (L218-19)	
15-17										

Tid	Må	09-nov	Ti	10-nov	On	11-nov	To	12-nov	Fr	13-nov
08-10	Kvantfysik (A)	Våglära/optik (A)	kap 16	Våglära/optik (A)	kap 16+32	Våglära/optik (A)	Våglära/optik (A)	kap 32+33	Kvantfysik (A)	
10-12	Våglära/optik kap 15+16	ÄFYA11 (L218)	Kvantfysik (A)	Kvantfysik (A)	Kvantfysik (A)	Våglära/optik (A)	Våglära/optik (A)	kap 33	Våglära/optik (A)	
13-15	Sl gp1-5 (L219)	ÄFYA11 (L218)	Övningar Optik&Våg (L218-19)	ÄFYA11 (L218)	Sl gp6-10 (L219)	Sl gp1-5 (L218)	Sl gp11-15 (L219)		Övningar Optik&Våg (L218-19)	
15-17										

Tid	Må	16-nov	Ti	17-nov	On	18-nov	To	19-nov	Fr	20-nov
08-10	Kvantfysik (A)	Våglära/optik (A)	kap 34	Kvantfysik (A)	Kvantfysik (A)	Våglära/optik (A)	kap 35	ÄFYA11 (L218)	Våglära/optik (A)	kap 36
10-12	Våglära/optik kap 34	Kvantfysik (A)	Kvantfysik (A)	Våglära/optik (A)	kap 34+35	Våglära/optik (A)	kap 36		Kvantfysik (A)	
13-15	Sl gp6-10 (L219)	Övningar Optik&Våg (L218-19)	Seminar/gen.gång (A) 12-13	Sl gp1-5 (L218) 13-15	Sl gp11-15 (L219) 13-15	Labbintroduktion (A) O2-O3, K1-K3			Övningar Optik&Våg (L218-19)	
15-17										

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Sound & Pressure



Sound as pressure waves

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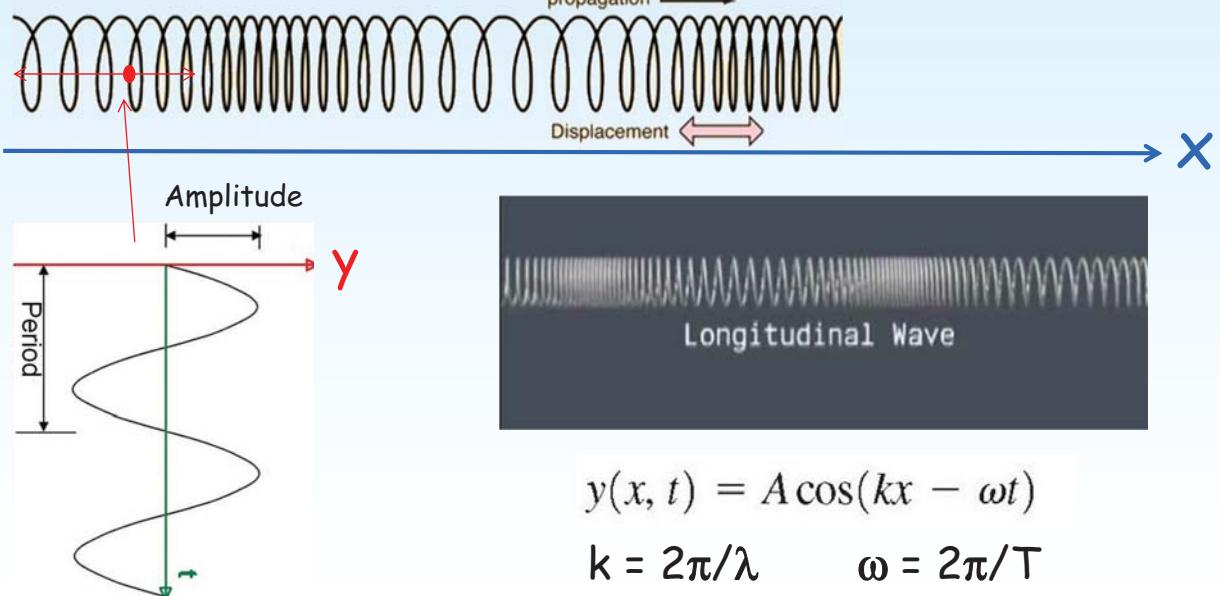


Sound & Pressure



Longitudinal sinusoidal wave

$$v = \omega / k$$



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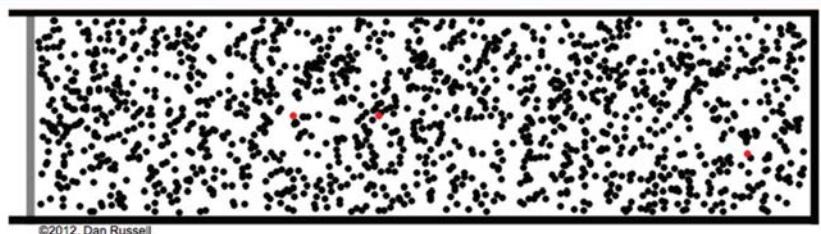
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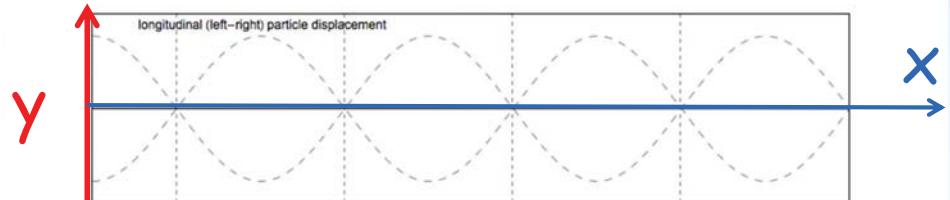
Sound & Pressure



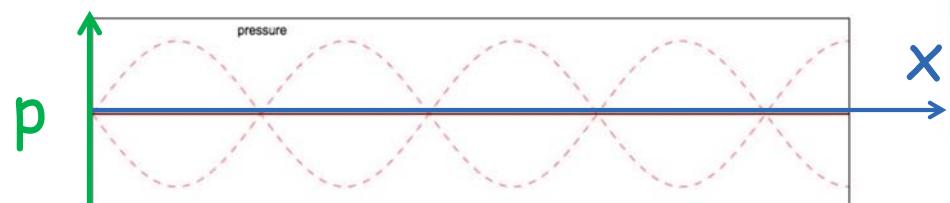
Piston moving
in and out:



Air molecule
movement:



Pressure:



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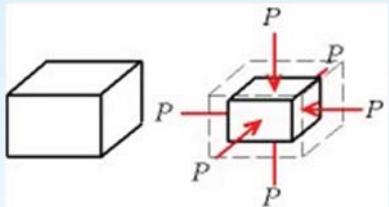
6



Sound & Pressure



Bulk modulus



The bulk modulus measures a medium's resistance to uniform compression:

$$B = -V \frac{\Delta p}{\Delta V} \quad \begin{array}{l} \xrightarrow{\text{Pressure change}} \\ \xrightarrow{\text{Volume change}} \end{array}$$

The change in pressure after a change of volume:

$$\Delta p = -B \frac{\Delta V}{V}$$

↑

Pressure increase: $\Delta p > 0$ and $\Delta V < 0$

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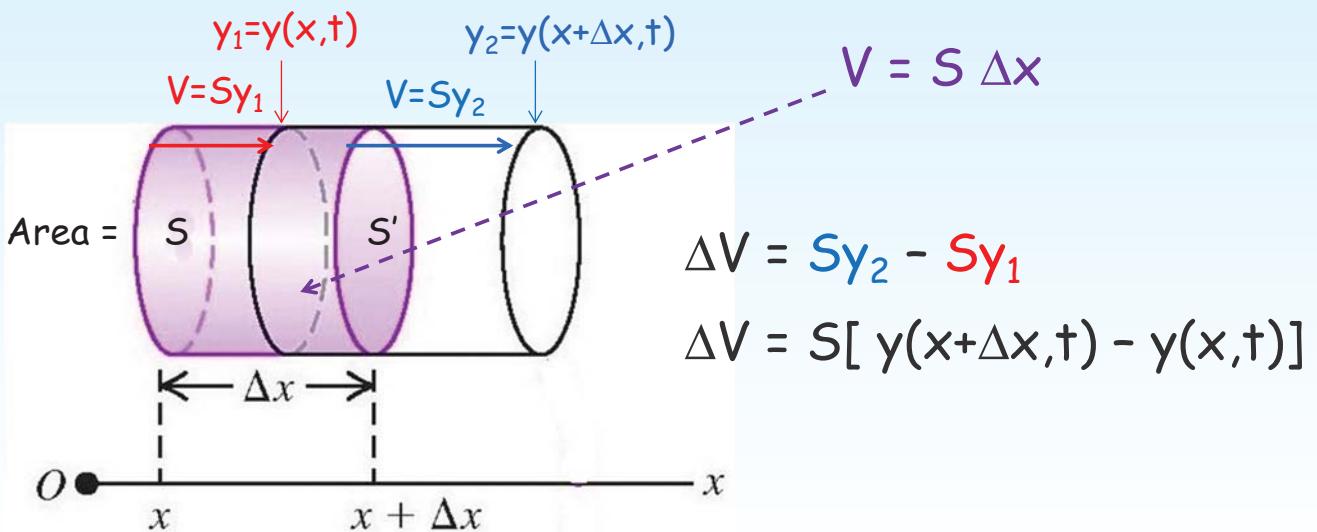
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Sound & Pressure



A soundwave is moving the area S to y_1 and the area S' to y_2 .



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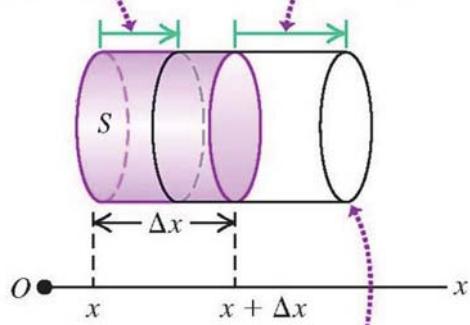
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Sound & Pressure



A sound wave
displaces the left end of the cylinder by $y_1 = y(x, t)$...
... and the right end by $y_2 = y(x + \Delta x, t)$.



The change in volume of the disturbed cylinder of fluid is $S(y_2 - y_1)$.

$$y(x, t) = A \cos(kx - \omega t)$$

$$\Delta V = S(y_2 - y_1) = S[y(x + \Delta x, t) - y(x, t)]$$

$$V = S \Delta x$$

$$\frac{dV}{V} = \lim_{\Delta x \rightarrow 0} \frac{S[y(x + \Delta x, t) - y(x, t)]}{S \Delta x} = \frac{\partial y(x, t)}{\partial x}$$

$$\Delta p = -B \frac{\Delta V}{V}$$

Let $p(x, t)$ be the instantaneous pressure fluctuation

$$p(x, t) = -B \frac{\partial y(x, t)}{\partial x}$$

$$p(x, t) = BkA \sin(kx - \omega t)$$

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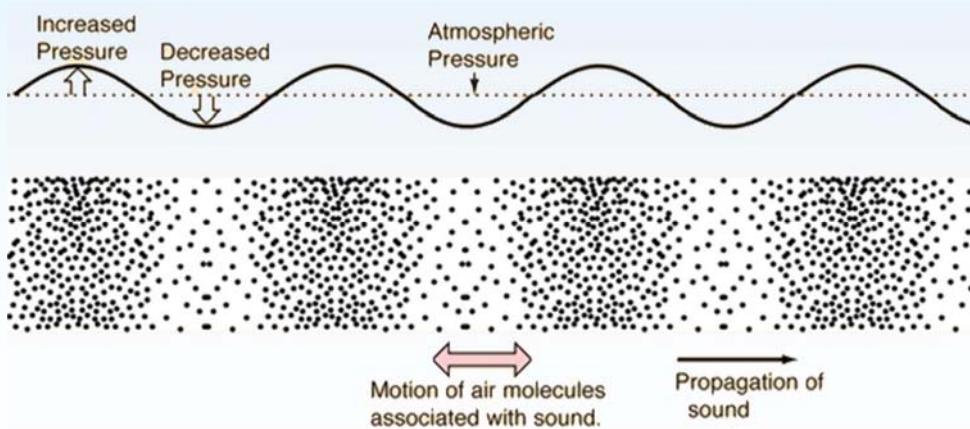
Sound & Pressure



$$p(x, t) = BkA \sin(kx - \omega t)$$

The pressure amplitude
The maximum pressure fluctuation

$$p_{\max} = BkA$$



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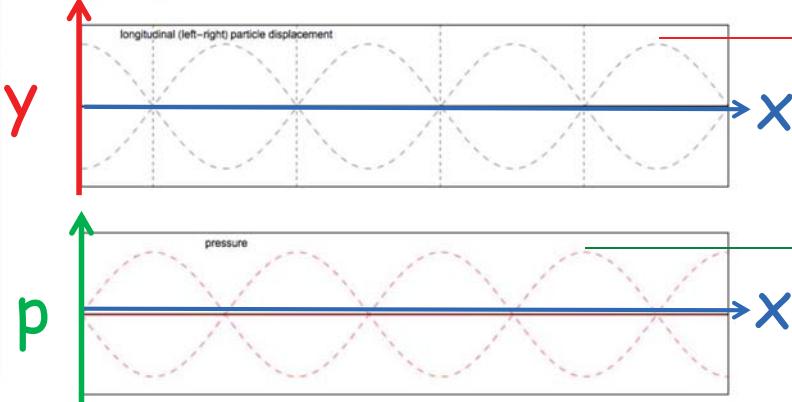
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Sound & Pressure



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$$y(x, t) = A \cos(kx - \omega t)$$

$$p(x, t) = BkA \sin(kx - \omega t)$$

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Sound & Pressure



Human hearing

Audible range: 20-20 kHz the human frequency range.

Loudness: Higher pressure amplitude \rightarrow Higher loudness
(at constant frequency)

Different frequency \rightarrow Different loudness
(at constant pressure amplitude)

Pitch: Higher frequency \rightarrow High pitch

Higher pressure amplitude \rightarrow Usually higher pitch

Timbre: Tone color or harmonic content.

Instruments with the same fundamental frequency can have different harmonic content.

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Sound & Problems



Problem solving

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Sound & Problems



In a sinusoidal sound wave of moderate loudness, the maximum pressure variations are about 3.0×10^{-2} Pa above and below atmospheric pressure. Find the corresponding maximum displacement if the frequency is 1000 Hz. In air at normal atmospheric pressure and density, the speed of sound is 344 m/s and the bulk modulus is 1.42×10^5 Pa.

$$k = \frac{\omega}{v} = \frac{2\pi f}{v} = \frac{(2\pi \text{ rad})(1000 \text{ Hz})}{344 \text{ m/s}} = 18.3 \text{ rad/m}$$

$$p_{\max} = BkA \rightarrow A = \frac{p_{\max}}{Bk} = \frac{3.0 \times 10^{-2} \text{ Pa}}{(1.42 \times 10^5 \text{ Pa})(18.3 \text{ rad/m})} = 1.2 \times 10^{-8} \text{ m}$$

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Sound - velocity



The velocity of sound in a liquid

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Sound - velocity



Kinematics

$$\text{Momentum: } \vec{p} = m\vec{v}$$

$$\text{Impuls: } \vec{J} = \int_{t_1}^{t_2} \sum \vec{F} dt$$

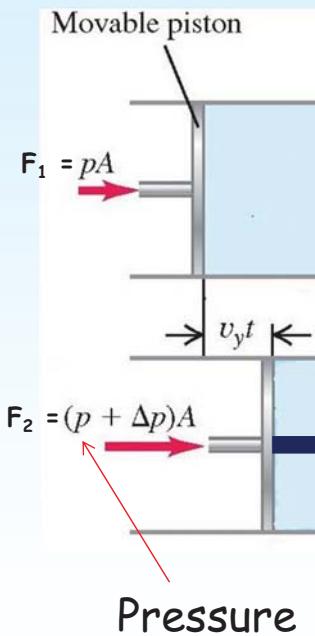
$$\text{The Momentum-Impuls theorem: } \vec{J} = \vec{p}_2 - \vec{p}_1$$

The impulse is equal to the change of momentum !

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Sound - velocity



Liquid

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$$\vec{J} = \int_{t_1}^{t_2} \sum \vec{F} dt$$

$$\vec{J} = (\vec{F}_2 - \vec{F}_1) t = \Delta p A t$$

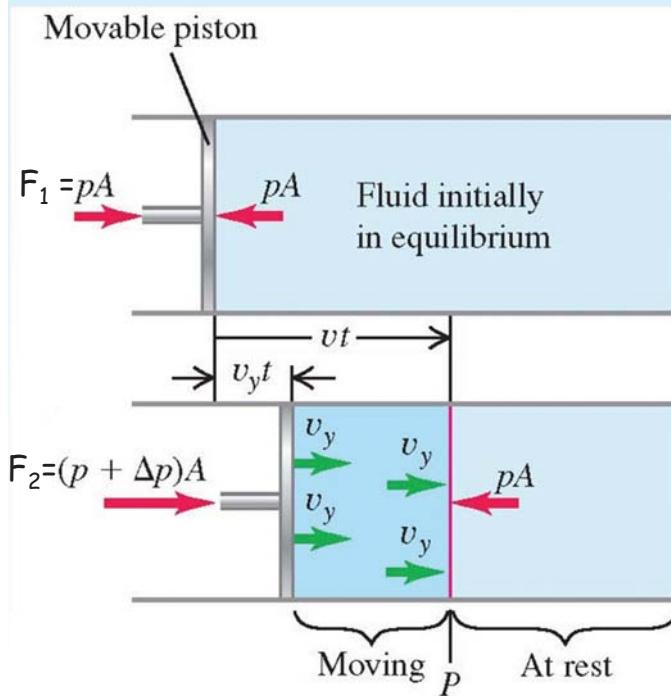
$$\vec{J} = \vec{p}_2 - \vec{p}_1$$

$$\vec{J} = \vec{P}_2 - \vec{P}_1 = \vec{P}_2 - 0 = m v_y$$

Momentum

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Sound - velocity



Sound in a liquid

Time = 0:

p : pressure in the liquid

A : area of the piston

F_1 : force on the piston

ρ : density of the liquid

Time = t :

v_y = velocity of the piston

v = velocity of the wave

$v_y t$ = distance the piston has moved

$v t$ = distance the wave has moved

Δp = increase of pressure

F_2 : force on the piston

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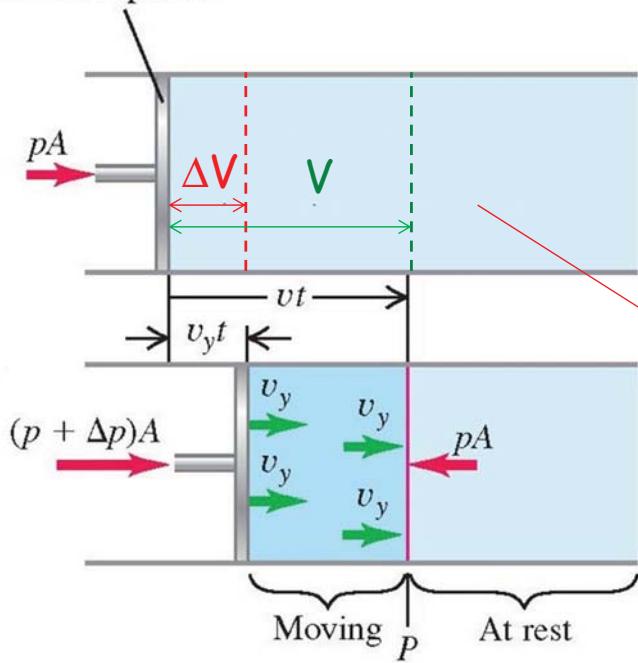
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Sound - velocity



Movable piston



The bulk modulus measures a medium's resistance to uniform compression:

$$B = -V \frac{\Delta p}{\Delta V} \quad \begin{array}{l} \text{Pressure change} \\ \text{Volume change} \end{array}$$

$$\Delta p = -B \frac{\Delta V}{V}$$

$$V = At$$

$$\Delta V = -Av_y t$$

Volume is decreasing

$$\Delta p = B \frac{v_y}{v}$$

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Sound - velocity

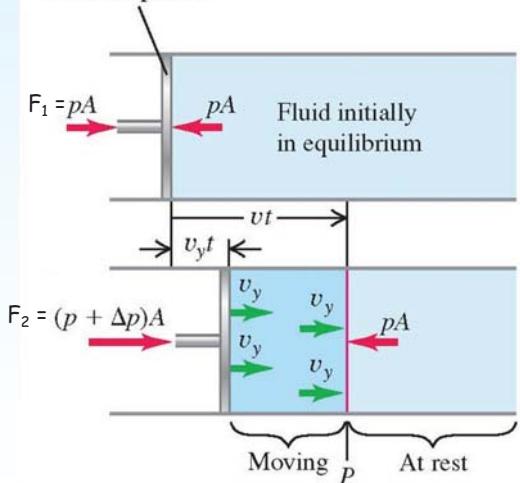


$$\vec{J} = \int_{t_1}^{t_2} \sum \vec{F} dt$$

$$\vec{J} = (\vec{F}_2 - \vec{F}_1) t = \Delta p A t = B A t v_y / v$$

$$\Delta p = B \frac{v_y}{v}$$

Movable piston



$$\vec{J} = \vec{p}_2 - \vec{p}_1 = \vec{p}_2$$

$$\vec{p} = m \vec{v}$$

$$\vec{J} = \vec{P}_2 = m v_y = \rho V v_y = \rho A v t v_y$$

$$\vec{J} = \vec{J}$$

$$\rho \frac{v_y}{v} A t = \rho v t A v_y$$

$$v = \sqrt{\frac{B}{\rho}} \quad (\text{speed of a longitudinal wave in a fluid})$$

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Sound - velocity



General:

$$v = \sqrt{\frac{\text{Restoring force returning the system to equilibrium}}{\text{Inertia resisting the return to equilibrium}}}$$

String:

$$v = \sqrt{\frac{F}{\mu}}$$

F: String tension
 μ : Mass per unit length

Liquid:

$$v = \sqrt{\frac{B}{\rho}}$$

B: Bulk modulus
 ρ : Density

Solid:

$$v = \sqrt{\frac{Y}{\rho}}$$

Y: Young's module
 ρ : Density

Gas:

$$v = \sqrt{\frac{B}{\rho}} = \sqrt{\frac{\gamma P}{\rho}} = \sqrt{\frac{\gamma RT}{M}}$$

γ : Adiabatic index
P: Pressure = nRT / V
 ρ : Density = m/V
R: Gas constant = 8.31 J/mol per K
T: Absolute temperature in K
M: Molar mass = m / n



Sound & Problems



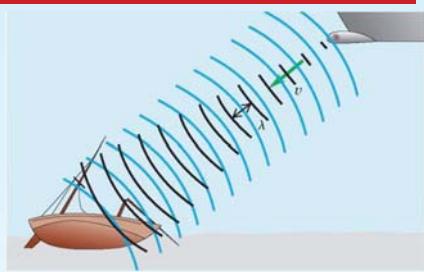
Problem solving



Sound & Problems



A ship uses a sonar system to locate underwater objects. Find the speed of sound waves in water and find the wavelength of a 262-Hz wave. $B = 2.18 \cdot 10^9 \text{ Pa}$ for water



$$v = \sqrt{\frac{B}{\rho}} = \sqrt{\frac{2.18 \times 10^9 \text{ Pa}}{1.00 \times 10^3 \text{ kg/m}^3}} = 1480 \text{ m/s}$$

$$\lambda = \frac{v}{f} = \frac{1480 \text{ m/s}}{262 \text{ s}^{-1}} = 5.65 \text{ m}$$

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Sound & Problems



Find the speed of sound in air at $T = 20^\circ\text{C}$, and find the range of wavelengths in air to which the human ear (which can hear frequencies in the range of 20–20,000 Hz) is sensitive. The mean molar mass for air (a mixture of mostly nitrogen and oxygen) is $M = 28.8 \times 10^{-3} \text{ kg/mol}$ and the ratio of heat capacities is $\gamma = 1.40$

$$T = 20^\circ\text{C} = 293 \text{ K}$$

$$v = \sqrt{\frac{\gamma RT}{M}} = \sqrt{\frac{(1.40)(8.314 \text{ J/mol} \cdot \text{K})(293 \text{ K})}{28.8 \times 10^{-3} \text{ kg/mol}}} = 344 \text{ m/s}$$

$$\lambda = v / f$$

$$\lambda = 344 / 20 = 17 \text{ m} \quad \text{for } f = 20 \text{ Hz}$$

$$\lambda = 344 / 20000 = 1.7 \text{ cm} \quad \text{for } f = 20 \text{ kHz}$$

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Sound – power & intensity



The power and intensity of sound

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Sound – power & intensity



The power in general:

$$P = \vec{F} \cdot \vec{v}$$

(instantaneous rate at which force \vec{F} does work on a particle)

Wave power (P):

The instantaneous rate at which energy is transferred along the wave.

Unit: W or J/s

$$P(x, t) = F_y(x, t)v_y(x, t) = -F \frac{\partial y(x, t)}{\partial x} \frac{\partial y(x, t)}{\partial t}$$

Wave intensity (I):

Average power per unit area through a surface perpendicular to the wave direction.

Unit: W/m²

$$I = P_{av} / \text{Area}$$

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Sound – power & intensity



The wave function:

$$y(x, t) = A \cos(kx - \omega t)$$

$$v_y(x, t) = \frac{\partial y(x, t)}{\partial t} = \omega A \sin(kx - \omega t)$$

The wave power:

$$P(x, t) = F_y(x, t)v_y(x, t)$$

The pressure function:

$$p(x, t) = BkA \sin(kx - \omega t)$$

Pressure is equal to force per unit area

The wave power per unit area:

$$\begin{aligned} P(x, t) &= p(x, t)v_y(x, t) = [BkA \sin(kx - \omega t)][\omega A \sin(kx - \omega t)] \\ &= B\omega k A^2 \sin^2(kx - \omega t) \end{aligned}$$

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Sound – power & intensity



The wave power per unit area:

$$P(x, t) = B\omega k A^2 \sin^2(kx - \omega t)$$

Intensity = Average wave power per unit area:

$$I = P_{av} / \text{Area} = \frac{1}{2} B \omega k A^2$$

Because the average of $\sin^2(x)$ is $\frac{1}{2}$

$$v = \omega / k$$

$$v = \sqrt{\frac{B}{\rho}}$$

$$k = \omega / \sqrt{\frac{B}{\rho}}$$

$$I = \frac{1}{2} \sqrt{\rho B} \omega^2 A^2 \quad (\text{intensity of a sinusoidal sound wave})$$

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Sound – power & intensity



$$I = \frac{1}{2} \sqrt{\rho B} \omega^2 A^2 \quad (\text{intensity of a sinusoidal sound wave})$$

The pressure amplitude
The maximum pressure fluctuation

$$p_{\max} = B k A$$

$$k = \omega / \sqrt{\frac{B}{\rho}}$$

$$p_{\max} = B A \omega / \sqrt{\frac{B}{\rho}} \quad \Rightarrow \quad A^2 \omega^2 = p_{\max}^2 / (\rho B)$$

$$I = \frac{p_{\max}^2}{2 \sqrt{\rho B}}$$

The intensity is proportional to the square of the pressure amplitude



Sound & Problems



Problem solving



Sound & Problems



Find the intensity of the sound wave with $p_{\max} = 3.0 \times 10^{-2}$ Pa. Assume the temperature is 20°C so that the density of air is $\rho = 1.20 \text{ kg/m}^3$ and the speed of sound is $v = 344 \text{ m/s}$.

$$I = \frac{p_{\max}^2}{2 \sqrt{\rho B}}$$

$$v = \sqrt{\frac{B}{\rho}} \quad \longrightarrow \quad v \rho = \sqrt{\rho B}$$

$$\begin{aligned} I &= \frac{p_{\max}^2}{2\rho v} = \frac{(3.0 \times 10^{-2} \text{ Pa})^2}{2(1.20 \text{ kg/m}^3)(344 \text{ m/s})} \\ &= 1.1 \times 10^{-6} \text{ J/(s} \cdot \text{m}^2) = 1.1 \times 10^{-6} \text{ W/m}^2 \end{aligned}$$

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Sound & Problems



What are the pressure and displacement amplitudes of a 20-Hz sound wave with the same intensity as the 1000-Hz sound wave of $p_{\max} = 3.0 \times 10^{-2}$ Pa, $\rho = 1.20 \text{ kg/m}^3$, $v = 344 \text{ m/s}$, $I = 1.1 \times 10^{-6} \text{ W/m}^2$

$$\begin{aligned} I &= \frac{1}{2} P_{\max}^2 / \sqrt{\rho B} \\ \text{since } \sqrt{\rho B} &\text{ is a constant } P_{\max} \text{ is the same if } I \text{ is the same.} \\ P_{\max} &= 3.0 \cdot 10^{-2} \text{ Pa} \end{aligned}$$

$$I = \frac{1}{2} \sqrt{\rho B} \omega^2 A^2 \quad (\text{intensity of a sinusoidal sound wave})$$

$$v = \sqrt{\frac{B}{\rho}} \quad \longrightarrow \quad v \rho = \sqrt{\rho B}$$

$$I = v \rho \omega^2 A^2 / 2 \quad \longrightarrow \quad A^2 = 2I / (v \rho \omega^2)$$

$$A^2 = 2 \times 1.1 \times 10^{-6} / (344 \times 1.20 \times (40\pi)^2) \quad A = 0.60 \mu\text{m}$$

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Sound & Problems



For an outdoor concert we want the sound intensity to be 1 W/m^2 at a distance of 20 m from the speaker array. If the sound intensity is uniform in all directions, what is the required acoustic power output of the array?

Intensity is average power per unit area

$$I = P_{\text{av}} / A_{\text{rea}}$$

The intensity through a sphere with radius r

$$I = \frac{P}{4\pi r^2}$$

The intensity through a hemisphere with radius r

$$I = \frac{P}{2\pi r^2}$$

$$P = 2 \pi r^2 I = 2.5 \text{ kW}$$

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Sound - Decibel



The decibel scale
of the intensity

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Sound - Decibel



Intensity in the unit of decibel (dB)

$$\beta = (10 \text{ dB}) \log \frac{I}{I_0} \quad (\text{definition of sound intensity level})$$

$I_0 = 10^{-12} \text{ W/m}^2$ is a reference intensity

It is roughly the threshold of human hearing

$\beta = 0 \text{ dB}$ for $I = I_0$

$\beta = 120 \text{ dB}$ for $I = 1 \text{ W/m}^2$



Sound - Decibel



Source or Description of Sound	Sound Intensity Level, β (dB)	Intensity, I (W/m^2)
Military jet aircraft 30 m away	140	10^2
Threshold of pain	120	1
Riveter	95	3.2×10^{-3}
Elevated train	90	10^{-3}
Busy street traffic	70	10^{-5}
Ordinary conversation	65	3.2×10^{-6}
Quiet automobile	50	10^{-7}
Quiet radio in home	40	10^{-8}
Average whisper	20	10^{-10}
Rustle of leaves	10	10^{-11}
Threshold of hearing at 1000 Hz	0	10^{-12}



Problem solving

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A 10-min exposure to 120-dB sound will temporarily shift your threshold of hearing at 1000 Hz from 0 dB up to 28 dB. Ten years of exposure to 92-dB sound will cause a *permanent* shift to 28 dB. What sound intensities correspond to 28 dB and 92 dB?

$$\beta = (10 \text{ dB}) \log \frac{I}{I_0}$$

$$I = I_0 10^{\beta/10} \quad I_0 = 10^{-12} \text{ W/m}^2$$

$$I_{28 \text{ dB}} = (10^{-12} \text{ W/m}^2) 10^{2.8} = 6.3 \times 10^{-10} \text{ W/m}^2$$

$$I_{92 \text{ dB}} = (10^{-12} \text{ W/m}^2) 10^{9.2} = 1.6 \times 10^{-3} \text{ W/m}^2$$

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Sound & Problems



Consider an idealized bird (treated as a point source) that emits constant sound power, with intensity obeying the inverse-square law
If you move twice the distance from the bird, by how many decibels does the sound intensity level drop?



$$\beta = (10 \text{ dB}) \log \frac{I}{I_0}$$

$$r_1 \quad \beta_1 \quad I_1 \quad r_2 = 2r_1 \quad \beta_2 \quad I_2$$

$$\begin{aligned} \beta_2 - \beta_1 &= (10 \text{ dB}) \left(\log \frac{I_2}{I_0} - \log \frac{I_1}{I_0} \right) \\ &= (10 \text{ dB}) [(\log I_2 - \log I_0) - (\log I_1 - \log I_0)] \\ &= (10 \text{ dB}) \log \frac{I_2}{I_1} \end{aligned}$$

$$\frac{I_1}{I_2} = \frac{r_2^2}{r_1^2} \quad (\text{inverse-square law for intensity})$$

$$\beta_2 - \beta_1 = (10 \text{ dB}) \log \frac{I_2}{I_1} = (10 \text{ dB}) \log \frac{1}{4} = -6.0 \text{ dB}$$

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Sound – Standing waves



Sound and standing waves

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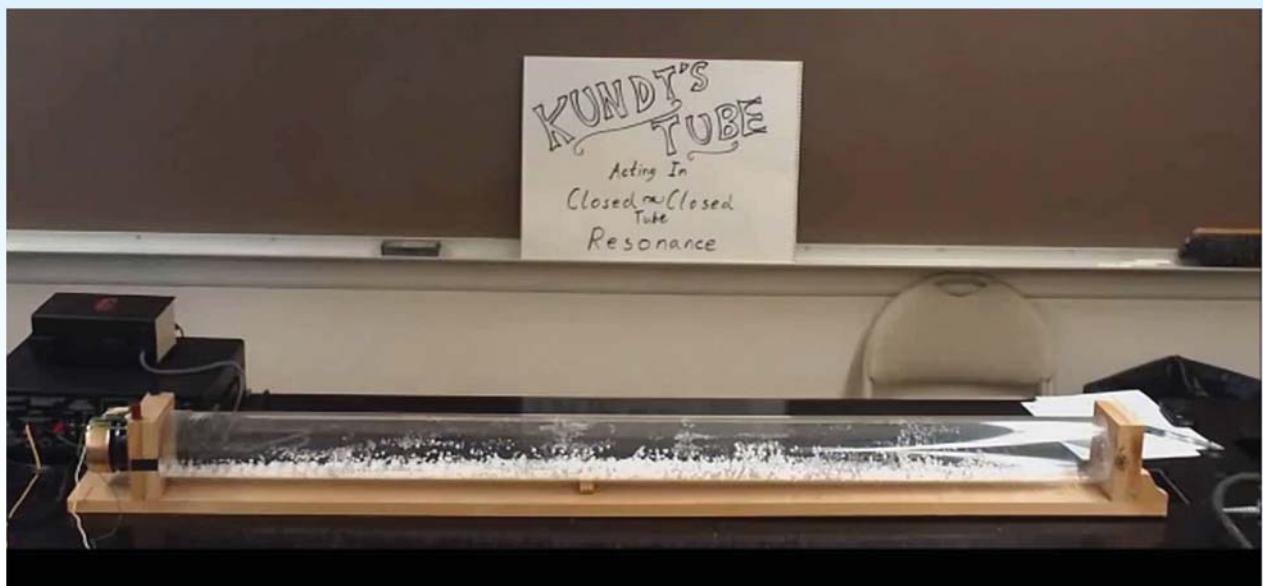
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Sound – Standing waves



Kundt's tube

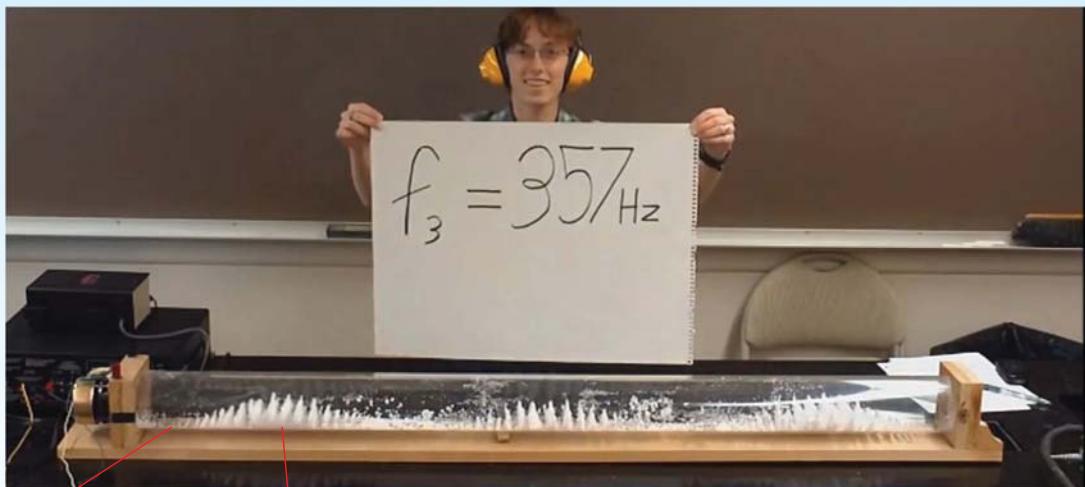


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Sound – Standing waves



$$\lambda = 95 \text{ cm}$$

$$v = \lambda f = 0.95 \times 357 = 339 \text{ m/s}$$

Displacement antinode
Maximum movement
Minimum pressure changes

Displacement node
Minimum movement
Maximum pressure changes

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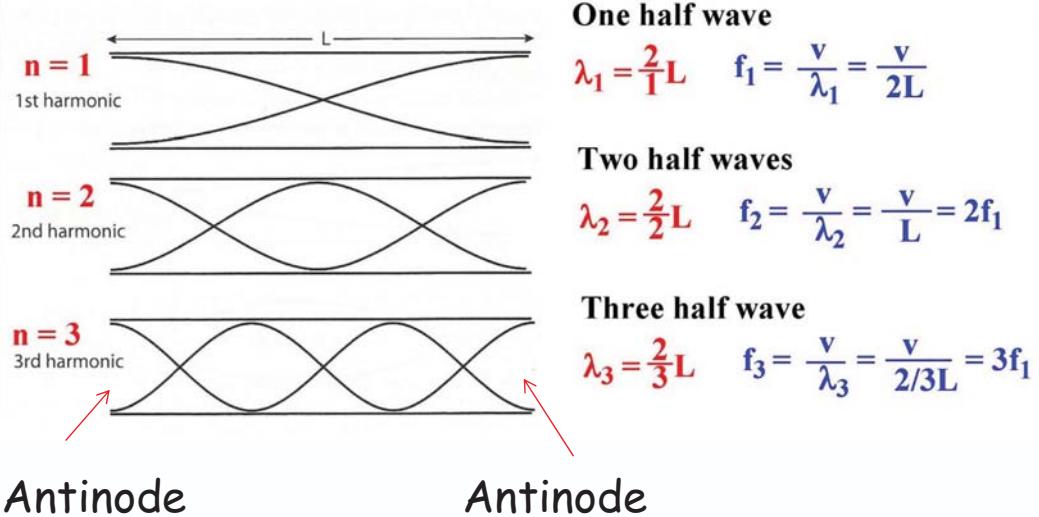


Sound – Standing waves



Standing wave in an open pipe

$$\lambda_n = \frac{2}{n}L \quad f_n = \frac{v}{\lambda_n} \quad \text{where the velocity (v) is the same for all } n$$



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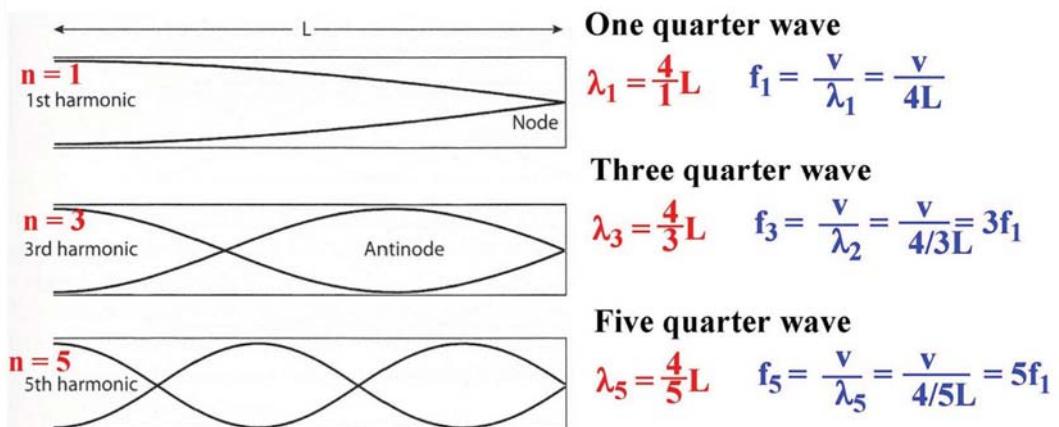
Sound – Standing waves



Standing wave in a closed pipe

$$\lambda_n = \frac{4}{n}L \quad f_n = \frac{v}{\lambda_n} \quad \text{where the velocity (v) is the same for all } n$$

Here the pressure is atmospheric giving displacement antinode (pressure node)



NOTE that $n = 2, 4, 6$ cannot happen in a closed pipe

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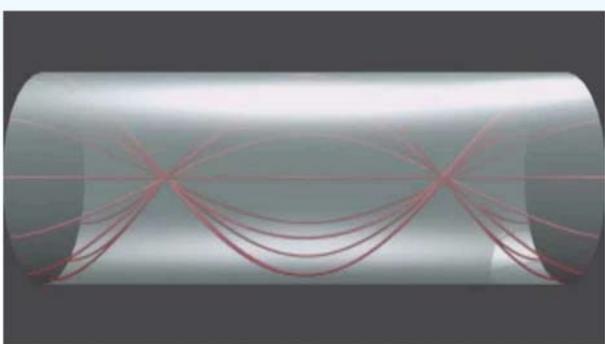
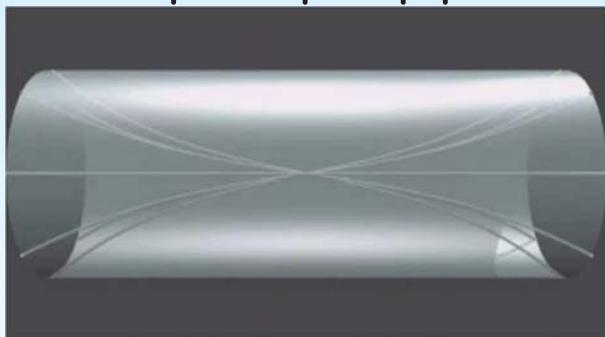
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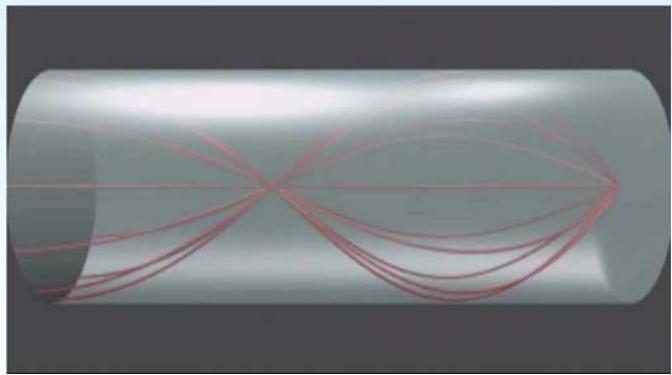
Sound – Standing waves



Open-open pipe



Open-closed pipe



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Sound – Standing waves

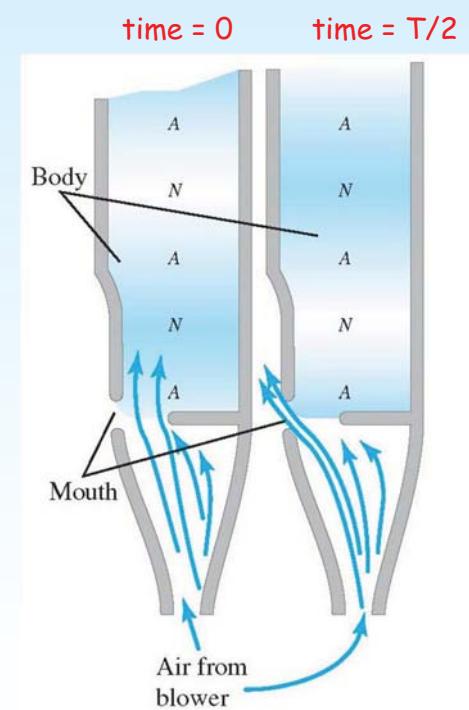


Organpipe: Airflow from below.

Standing wave: If the airspeed and pipelengths are chosen correctly.

Mouth: Pipe is open at the bottom and gives a pressure node (displacement antinode).

Airflow: Depending on time the air flow will either go into the pipe or out through the mouth.



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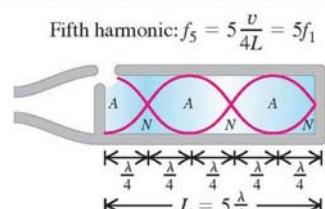
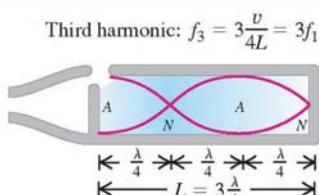
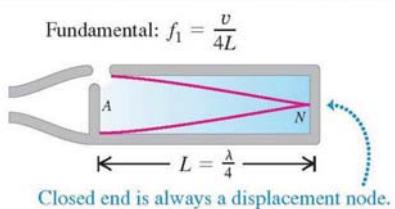
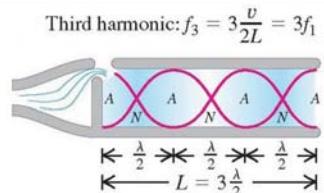
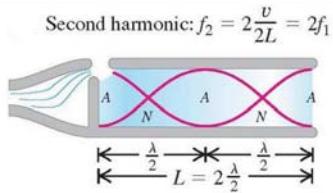
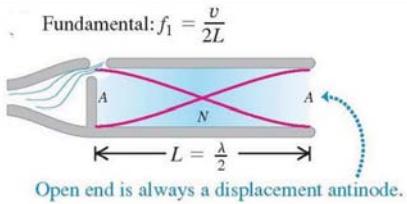
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Sound – Standing waves



The pipe can be open-open or open-closed



Remember: The distance between two nodes is $\lambda/2$

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Sound & Problems



Problem solving

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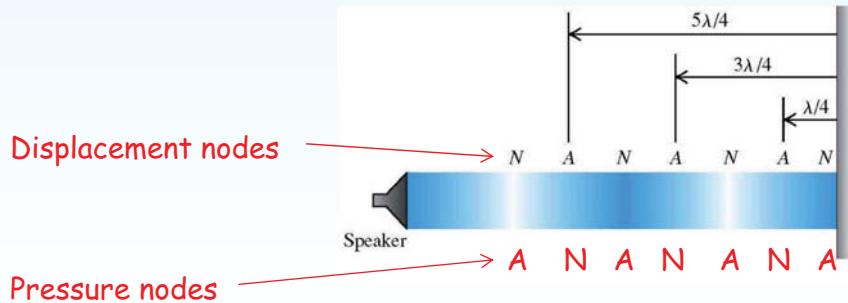


Sound & Problems



A directional loudspeaker directs a sound wave of wavelength λ at a wall
At what distances from the wall could you stand and hear no sound at all?

Your ear detects pressure variations in the air; you will therefore hear no sound if your ear is at a *pressure node*, which is a displacement antinode. The wall is at a displacement node.



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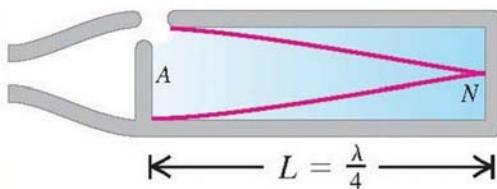


Sound & Problems



On a day when the speed of sound is 345 m/s, the fundamental frequency of a particular stopped organ pipe is 220 Hz. How long is this pipe?

$$\text{Fundamental: } f_1 = \frac{v}{4L}$$



$$L_{\text{stopped}} = \frac{v}{4f_1} = \frac{345 \text{ m/s}}{4(220 \text{ s}^{-1})} = 0.392 \text{ m}$$

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Sound & Problems



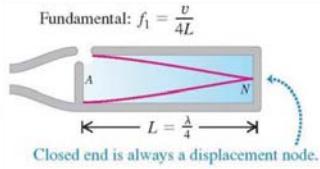
The second *overtone* of this pipe has the same wavelength as the third *harmonic* of an *open* pipe. How long is the open pipe?

$$v = 345 \text{ m/s}$$

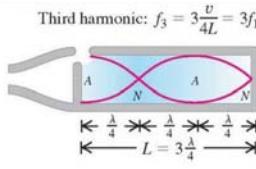
$$f_1 = 220 \text{ Hz}$$

$$L_{\text{stopped}} = 0.932 \text{ m}$$

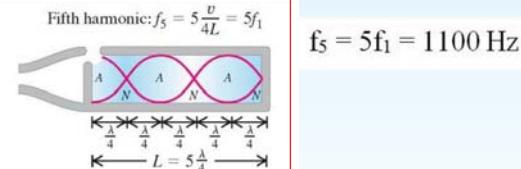
Fundamental frequency



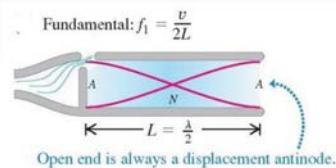
First overtone



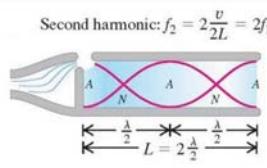
Second overtone



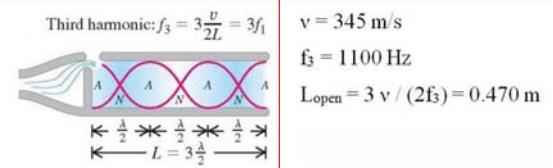
Fundamental



Second harmonic



Third harmonic



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Sound – Resonance



Sound and resonance

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Sound & Problems



Resonance



Many mechanical systems have **normal mode frequencies of oscillation**. In these modes every particle in the system oscillates with simple harmonic oscillation.

If an **outside drivingforce** is applied that varies with a **normal mode frequency** then the system is in **resonance** and the amplitude of the oscillations can increase.

In this case the drivingforce is **continuously adding energy** to the system.

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Sound & Problems



Problem solving

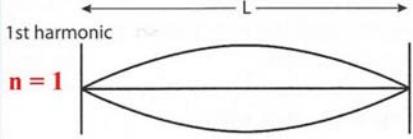
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Sound & Problems



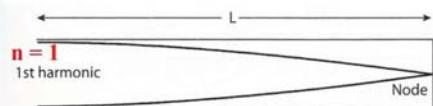
A stopped organ pipe is sounded near a guitar, causing one of the strings to vibrate with large amplitude. We vary the string tension until we find the maximum amplitude. The string is 80% as long as the pipe. If both pipe and string vibrate at their fundamental frequency, calculate the ratio of the wave speed on the string to the speed of sound in air.



One half wave

$$\lambda_1 = \frac{2}{1}L \quad f_1 = \frac{v}{\lambda_1} = \frac{v}{2L}$$

$$v_{\text{string}} = 2L_{\text{string}} f_{\text{string}}$$



One quarter wave

$$\lambda_1 = \frac{4}{1}L \quad f_1 = \frac{v}{\lambda_1} = \frac{v}{4L}$$

$$v_{\text{pipe}} = 4L_{\text{pipe}} f_{\text{pipe}}$$

$$v_{\text{string}} / v_{\text{pipe}} = 2L_{\text{string}} f_{\text{string}} / 4L_{\text{pipe}} f_{\text{pipe}}$$

$$f_{\text{string}} = f_{\text{pipe}}$$

$$L_{\text{string}} = 0.80 L_{\text{pipe}}$$

$$v_{\text{string}} / v_{\text{pipe}} = 2 \cdot 0.80 L_{\text{pipe}} f_{\text{pipe}} / 4L_{\text{pipe}} f_{\text{pipe}} = 0.40$$

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Sound – Interference



Sound and interference

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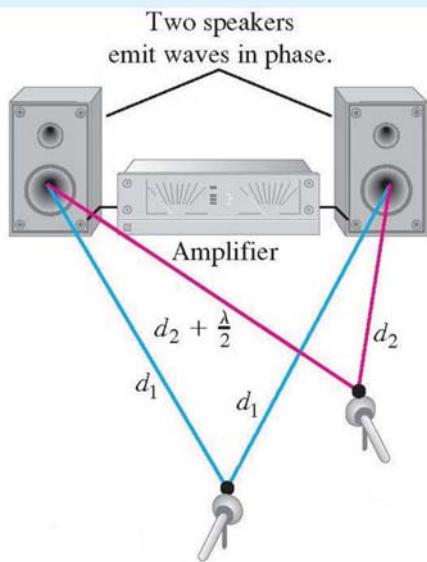
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Sound – Interference



Interference



Two waves that arrives at a point where the distance is different with $n\lambda$ ($n=0,1,2,3\dots$) undergo **constructive interference** and have a doubled amplitude.

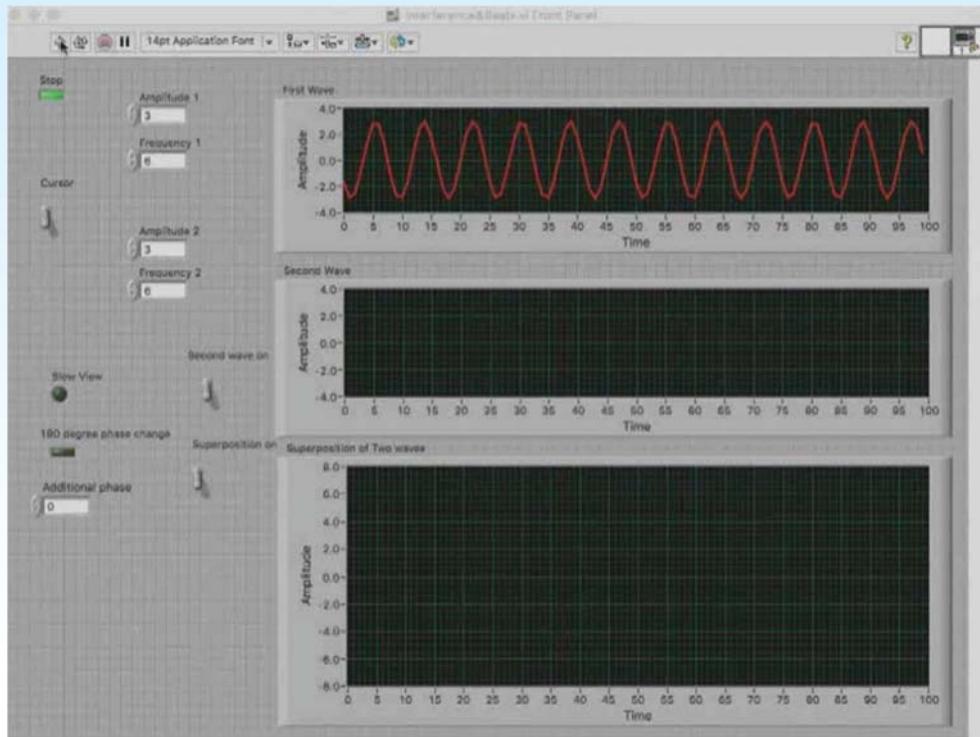
Two waves that arrives at a point where the distance is different with $n\lambda/2$ ($n=1,3,5\dots$) undergo **destructive interference** and have a zero amplitude.

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Sound – Interference



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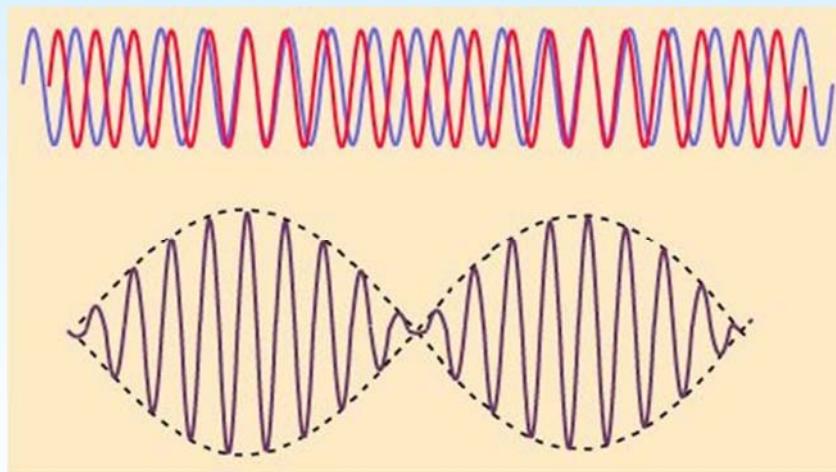
Sound – Interference



BEAT: If two sound waves with slightly different frequencies are added up they give a sound that is going up and down in intensity.

Two waves with
different
frequency

Their
superposition



This pulsating sound is only heard if the difference in frequency is < 7 Hz

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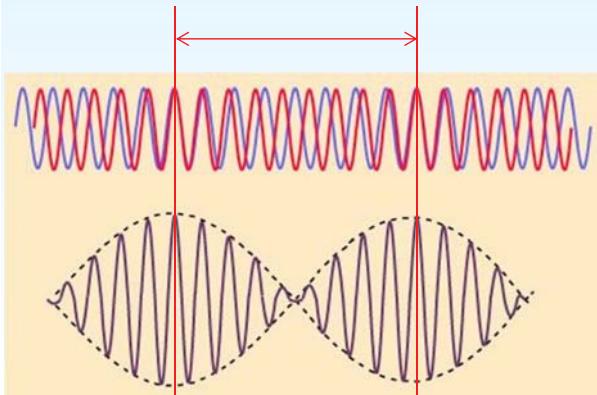
Sound – Interference



What is the frequency of the beat ?

$$T_{\text{beat}} = 9T_{\text{red}} = 8T_{\text{blue}}$$

$$T_{\text{beat}} = nT_a = (n-1)T_b$$



$$f_{\text{beat}} = f_a/n = f_b/(n-1)$$

$$f_{\text{beat}} = f_b/(n-1)$$

$$n = 1 + f_b/f_{\text{beat}} = (f_b + f_{\text{beat}}) / f_{\text{beat}}$$

$$f_{\text{beat}} = f_a/n$$

$$f_{\text{beat}} = f_a / [(f_b + f_{\text{beat}}) / f_{\text{beat}}] = f_a f_{\text{beat}} / (f_b + f_{\text{beat}})$$

$$1 = f_a / (f_b + f_{\text{beat}})$$

$$f_{\text{beat}} = f_a - f_b$$

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Sound – Doppler effect



The Doppler effect

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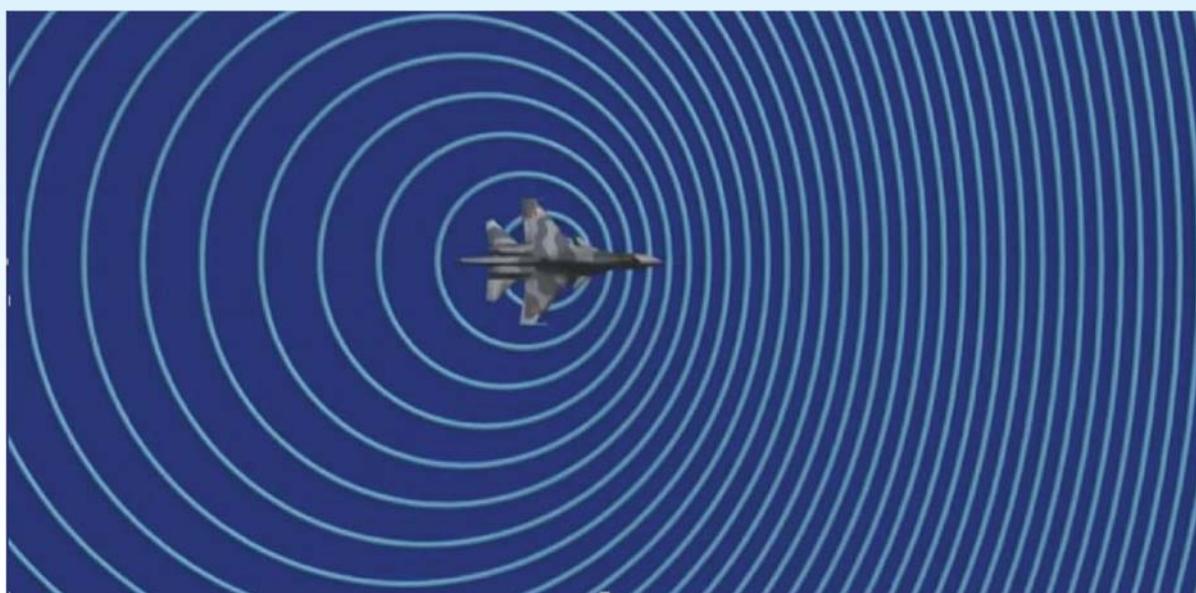
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Sound – Doppler effect



Doppler effect



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Sound – Doppler effect

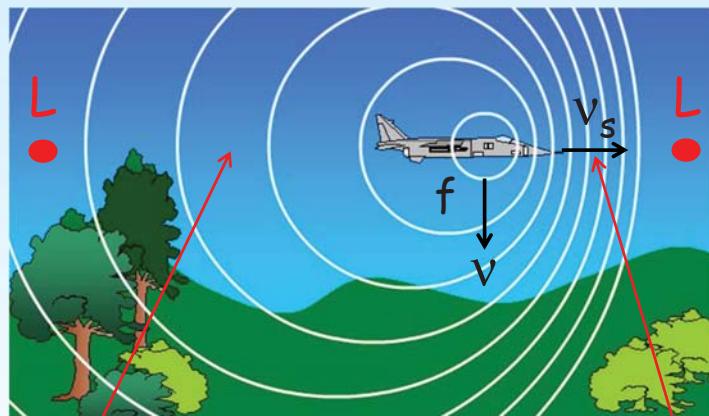


The time for a sound wave to reach a listener (L) gets longer if the source (S) is moving away.

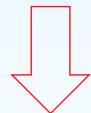


λ_{behind} longer

$$\lambda_{\text{behind}} = \frac{v + v_s}{f}$$



The time for a sound wave to reach a listener (L) gets shorter if the source is moving closer.



$\lambda_{\text{in front}}$ shorter

$$\lambda_{\text{in front}} = \frac{v - v_s}{f}$$

$$v = \lambda / T = \lambda f$$

$$\lambda = v / f$$

$$f = v / \lambda$$

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Sound – Doppler effect

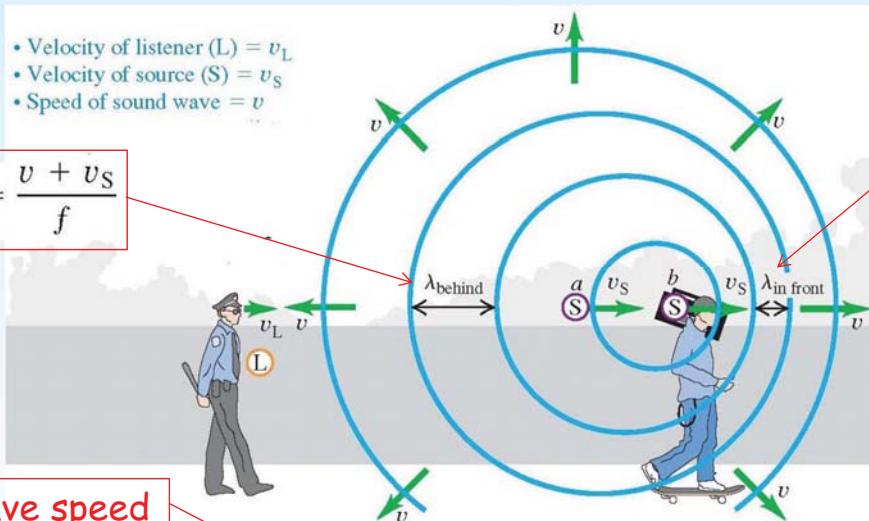


What if the listener is also moving ?

- Velocity of listener (L) = v_L
- Velocity of source (S) = v_s
- Speed of sound wave = v

$$\lambda_{\text{behind}} = \frac{v + v_s}{f}$$

$$\lambda_{\text{in front}} = \frac{v - v_s}{f}$$



The wave speed relative to L is

$$v + v_L$$

$$f_L = \frac{v + v_L}{\lambda_{\text{behind}}} = \frac{v + v_L}{(v + v_s)/f} = \frac{v + v_L}{v + v_s} f$$

change in frequency

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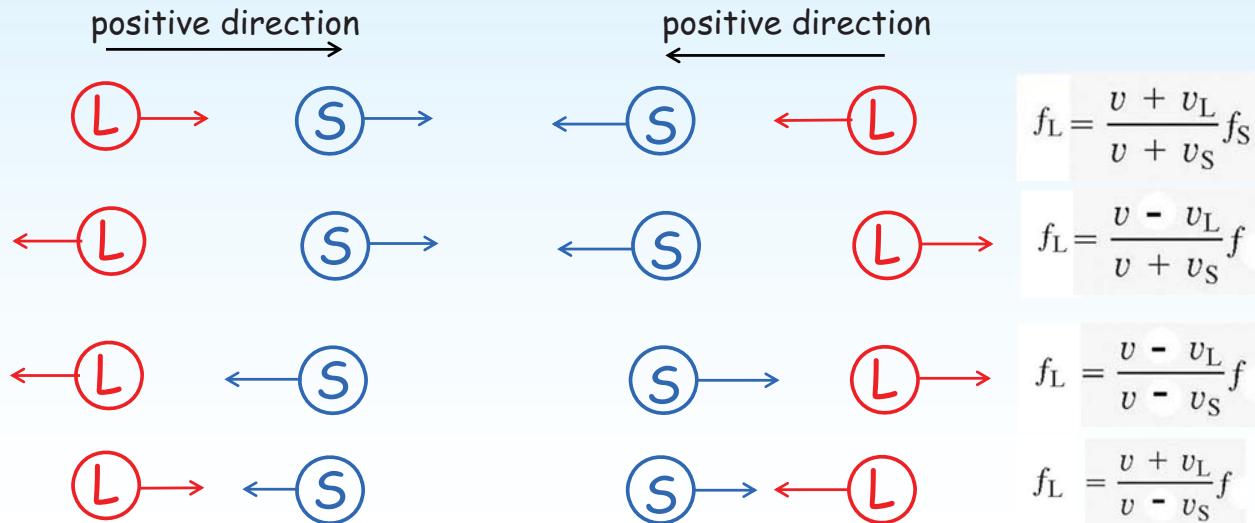


Sound – Doppler effect



$$f_L = \frac{v + v_L}{v + v_S} f_S$$

always works if the positive direction is defined as going from the listener to the source.



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Sound – Doppler effect



Electromagnetic waves such as light also have a Doppler shift. It can be calculated using the theory of relativity:

$$f_O = \sqrt{\frac{c - v}{c + v}} f_S$$

f_S = the frequency of the source

f_O = the frequency detected by an observer

c = the speed of light

v = the relative velocity of the source with respect to the observer

v is positive if the observer and the source is moving apart

v is negative if the observer and the source is moving towards each other

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Sound & Problems



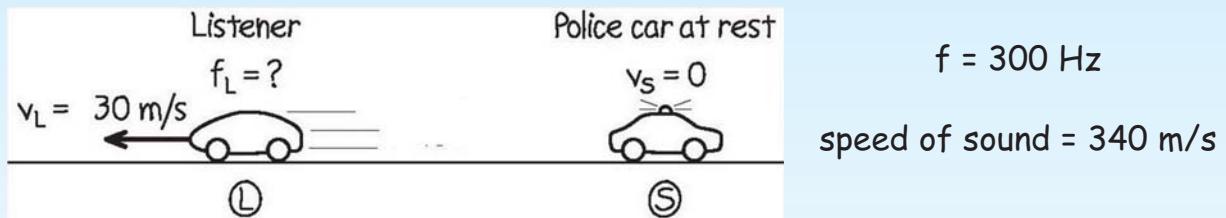
Problem solving

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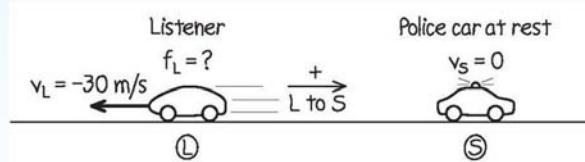
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Sound & Problems



What frequency does the listener hear ?



$$f_L = \frac{v + v_L}{v + v_S} f = \frac{340 \text{ m/s} + (-30 \text{ m/s})}{340 \text{ m/s}} (300 \text{ Hz}) = 274 \text{ Hz}$$

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Sound – shockwave



Shockwave

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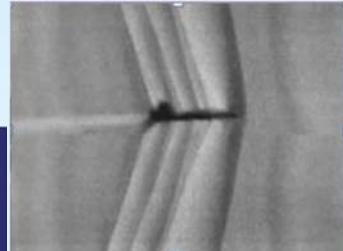
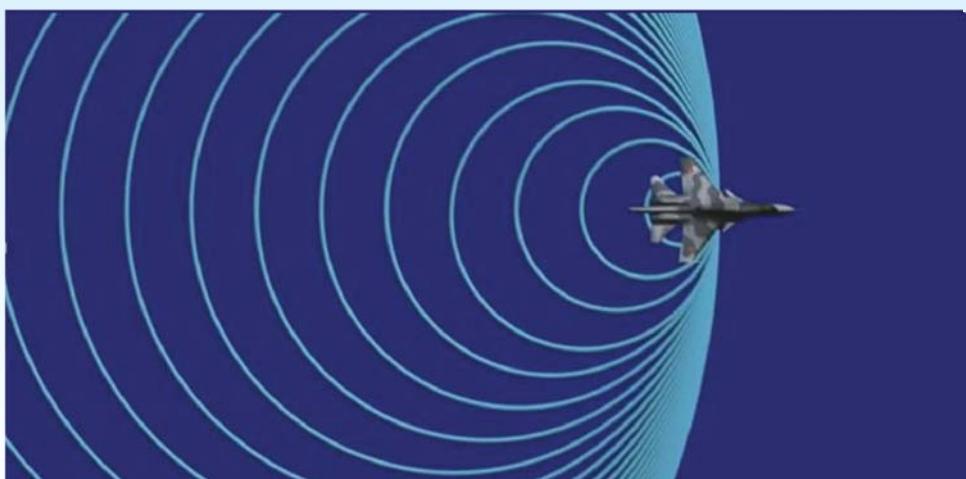
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Sound – shockwave



Shock waves



$$\lambda_{\text{in front}} = \frac{v - v_s}{f}$$

v: Speed of sound
 v_s : Speed of the plane

$v_s > v$ Shockwave is created (not only when $v_s = v$)

$v_s > v$ No sound in front of the plane

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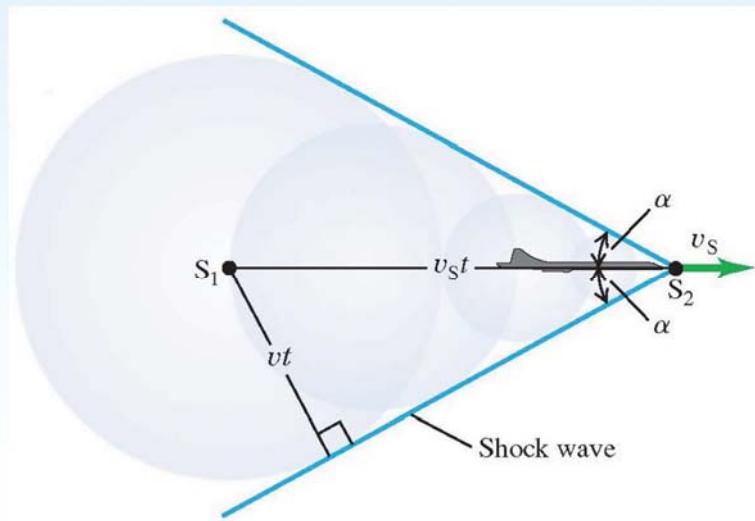


Sound



A conical shock wave is produced if a plane flies faster than the speed of sound.

A series of circular wave crests from the plane interfere constructively along a line that is given by an angle α .



v : Speed of sound
 v_s : Speed of the plane

Speed of the plane in Mach number:

$$N_M = v_s/v$$

$$\sin \alpha = \frac{vt}{v_s t} = \frac{v}{v_s}$$

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Sound & Problems



Problem solving

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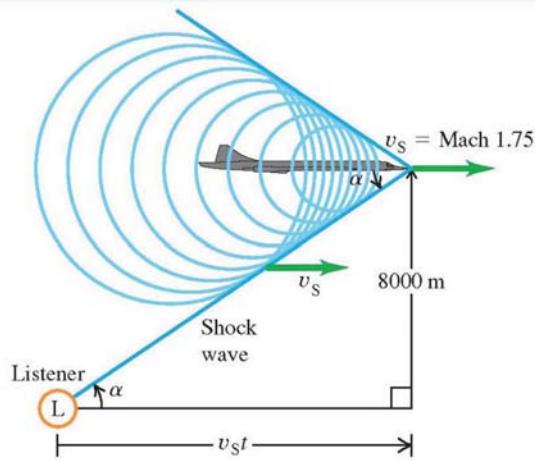
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Sound & Problems



An airplane is flying at Mach 1.75 at an altitude of 8000 m, where the speed of sound is 320 m/s. How long after the plane passes directly overhead will you hear the sonic boom?



$$N_M = v_s/v = 1.75$$

$$v_s = (1.75)(320 \text{ m/s}) = 560 \text{ m/s}$$

$$\sin \alpha = v / v_s = 1 / N_M = 1 / 1.75$$

$$\alpha = \arcsin \frac{1}{1.75} = 34.8^\circ$$

$$\tan \alpha = \frac{8000 \text{ m}}{v_s t}$$

$$t = \frac{8000 \text{ m}}{(560 \text{ m/s})(\tan 34.8^\circ)} = 20.5 \text{ s}$$